

# *Oscillating Quintessence*

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# Quintessence (based on FLRW)

Quintessence: dynamical scalar field  $\phi$

$$\text{Action : } S = \int dx^4 \sqrt{|g|} \left[ \frac{1}{2} g^{\mu\nu} (\partial_\mu \phi)(\partial_\nu \phi) - V(\phi) \right]$$

$$\text{Field equation: } \frac{\partial^2 \phi}{\partial t^2} + 3H \frac{\partial \phi}{\partial t} - \frac{1}{a^2} \nabla^2 \phi + V'(\phi) = 0$$

$$\text{energy density } \rho_\phi = \frac{1}{2} \left( \frac{\partial \phi}{\partial t} \right)^2 + \frac{1}{2a^2} (\nabla \phi)^2 + V(\phi)$$

$$\text{and pressure : } p_\phi = \frac{1}{2} \left( \frac{\partial \phi}{\partial t} \right)^2 - \frac{1}{6a^2} (\nabla \phi)^2 - V(\phi)$$

Slow evolution and weak spatial dependence

→  $V(\phi)$  dominates →  $w_\phi \sim -1$  → Acceleration



**How to achieve it (naturally) ?**

# Quintessence (FLRW)

Quintessence: dynamical scalar field  $\phi$

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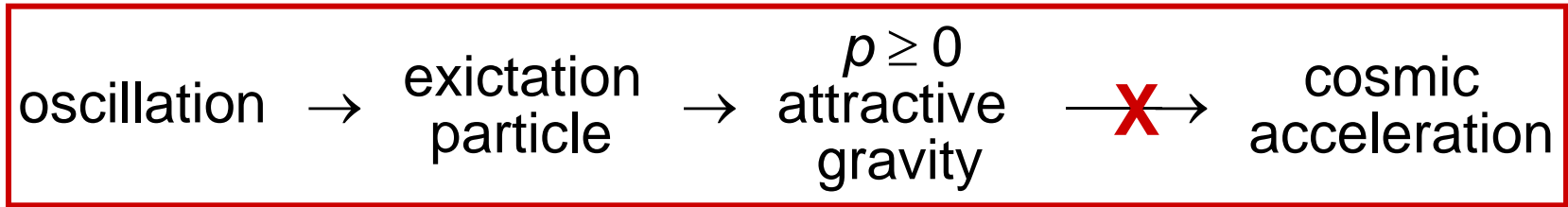
$$\text{and pressure :} \quad p_\phi = \frac{1}{2} \left( \frac{\partial \phi}{\partial t} \right)^2 - \frac{1}{6a^2} (\nabla \phi)^2 - V(\phi)$$

$$\text{Assume } \phi = \phi(t) \Rightarrow \rho_\phi = K + V, \quad p_\phi = K - V, \quad \text{where } K \equiv \frac{1}{2} \dot{\phi}^2$$

$$\Rightarrow w_\phi = \frac{p_\phi}{\rho_\phi} = \frac{(K - V)}{(K + V)}$$

# Quintessence: oscillation —?→ repulsive gravity

straightforward naïve thinking:



(example)  $V(\phi) = \frac{1}{2} m^2 \phi^2$  :

an oscillating mode :  $\phi \sim \phi_0 \cos(kx \pm \omega_k t + \theta)$  ,  $\omega_k^2 = \frac{k^2}{a^2} + m^2$

$$\langle \rho_\phi \rangle_t \cong \frac{1}{2} \left( \frac{k^2}{a^2} + m^2 \right) \phi_0^2 \quad , \quad \langle p_\phi \rangle_t \cong \frac{1}{6} \frac{k^2}{a^2} \phi_0^2 > 0$$

$$w_\phi = \frac{\langle p_\phi \rangle_t}{\langle \rho_\phi \rangle_t} \cong \begin{cases} \frac{1}{3} , & \text{for } \frac{k^2}{a^2} \gg m^2 \text{ (large - } k \text{ mode) } \sim \text{radiation} \\ 0 , & \text{for } \frac{k^2}{a^2} \ll m^2 \text{ (small - } k \text{ mode) } \sim \text{NR matter} \end{cases}$$

# Quintessence: a slowly evolving mode

For  $V(\phi) = \frac{1}{2} m^2 \phi^2$  :

solution:  $\phi \sim \phi_0 \cos(kx \pm \omega_k t + \theta)$ ,  $\omega_k^2 = \frac{k^2}{a^2} + m^2$

- weak spatial dependence:  $k \sim 0$

slow evolution:  $\omega_{k \sim 0} \leq H_0$

→  $m \leq H_0 \sim 10^{-48} \text{ GeV}$  (extremely small !!)

- $V(\phi) = \frac{1}{2} m^2 \phi^2 \sim \rho_\phi \sim \rho_c \equiv \frac{3H_0^2}{8\pi G}$  [V( $\phi$ ) dominates.]

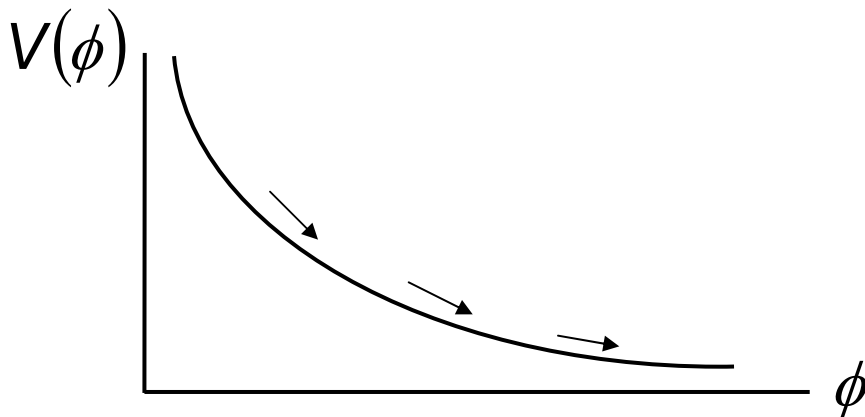
→  $\phi \geq G^{-1/2} \sim \text{Planck energy} \sim 10^{19} \text{ GeV}$  (extremely large !!)

# Quintessence: oscillation —?→ repulsive gravity



It seems that **oscillators** for repulsive gravity (dark energy).

Accordingly, instead of oscillation-like modes, most quintessence models involve **“slowly” moving modes** such as **running away behavior**:



(example) **tracker quintessence**

- Power-law :  $V(\phi) = \frac{M^{4+n}}{\phi^n}$
- Exponential :  $V(\phi) = V_0 e^{-\phi/M}$

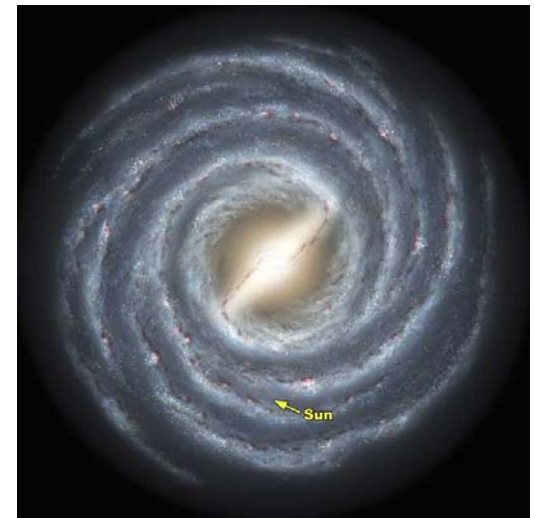
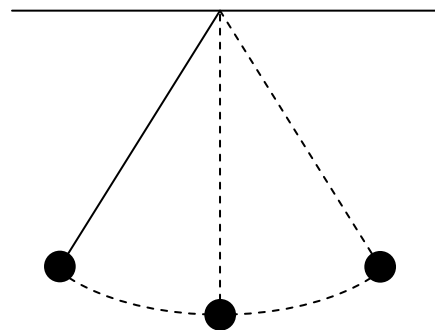
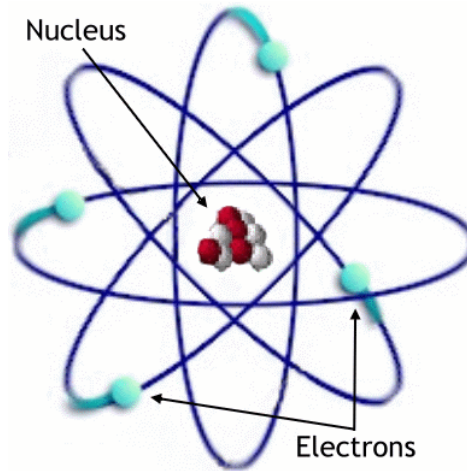
$\phi \rightarrow \infty$  ( $V \rightarrow 0$ ) **asymptotically.**

# Quintessence: oscillation —?→ repulsive gravity

**?** oscillation —~~X~~→ repulsive gravity (dark energy) !?

If so, pity..... (oscillation-like behavior so frequent/familiar/friendly)

(for example)



**Quintessence:** oscillation —?→ repulsive gravity

In this talk, show

# *Oscillating Quintessence*



**negative pressure**  
**repulsive gravity**  
dark energy

# Oscillating Quintessence

Action:  $S = \int dx^4 \sqrt{|g|} [K - V]$

Field equation:  $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

[ignoring spatial dependence  
ignoring spatial curvature ( $k/a^2$ )]

power-law potential:  $V(\phi) = M^{4-n} |\phi|^n$ ,  $n > 0$ ,  $M > 0$

(regularization for  $0 < n \leq 1$ :  $V(\phi) = \frac{M^{4-n} |\phi|^{1+n}}{|\phi| + \varepsilon}$ )

For oscillation, consider

**time-averaged quantities** (averaged over a period  $T$ ):  $\langle \rho_\phi \rangle, \langle p_\phi \rangle, \langle K \rangle, \langle V \rangle \dots$

$\Rightarrow \langle K \rangle = \frac{n}{2} \langle V \rangle \Rightarrow w_{OQ} = \frac{n-2}{n+2}$

(simply determined by  $n$ )

$n$	4	2	1	0.1	$0^+$
$w_{OQ}$	1/3	0	-1/3	-0.9	$-1^+$

$\Rightarrow 0 < n < 1 \Leftrightarrow -1 < w_{OQ} < -1/3$  (repulsive gravity)

# Oscillating Quintessence

(other quantities)

$$\rho_{\text{OQ}} \equiv \langle \rho_\phi \rangle = \rho_c \Omega_{\text{DE}} (1+z)^{3(1+w_{\text{OQ}})} = V_{\text{max}} = M^{4-n} |\phi|_{\text{max}}^n$$

$$\Omega_{\text{DE}} = \rho_{\text{OQ}}(0) / \rho_c = M^{4-n} |\phi|_{\text{max}}^n (0)$$

$$\text{period } T = \frac{\sqrt{8\pi}}{n} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} M^{1-4/n} \rho_{\text{OQ}}^{1/n-1/2}$$

$$HT = \frac{8\pi}{\sqrt{3n}} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} M_{\text{pl}}^{-1} M^{1-4/n} \rho^{1/2} \rho_{\text{OQ}}^{1/n-1/2} = \frac{8\pi}{\sqrt{3n}} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} \sqrt{\frac{\rho}{\rho_{\text{OQ}}}} \frac{|\phi|_{\text{max}}}{M_{\text{pl}}}$$

$$H_0 T_0 = \frac{\sqrt{3}}{n} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} \left(\frac{3}{8\pi}\right)^{4/n-1} \Omega_{\text{OQ}}^{1/n-1/2} \left(\frac{M}{M_{\text{pl}}}\right)^{1-2/n} \left(\frac{M}{H_0}\right)^{-2/n}$$

$$= \frac{8\pi}{\sqrt{3n}} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} \Omega_{\text{OQ}}^{-1/2} \frac{|\phi|_{\text{max}}(0)}{M_{\text{pl}}}$$

# Oscillating Quintessence

- Constraints from observations:  $w_{OQ} < -0.9 (1\sigma)$  (for a const.  $w_{DE}$ )  
 $\Omega_{DE} \cong 0.75$

$$\Rightarrow \begin{cases} 0 < n < 0.1 \\ V_{\max}(0) = M^{4-n} |\phi|_{\max}^n \cong 0.75 \rho_c \cong 3 \times 10^{-11} \text{eV}^4 \end{cases} \quad \left( \text{rem } w_{OQ} = \frac{n-2}{n+2} \right)$$

- Condition for oscillation:  $HT \ll 1$  from now back to some early time

$$H_0 T_0 \ll 1 \Rightarrow \begin{cases} M \gg M_{\min} \equiv \left[ \frac{8\pi}{\sqrt{3n}} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} \Omega_{OQ}^{1/n - 1/2} \rho^{1/n} M_{pl}^{-1} \right]^{n/4-n} \\ |\phi|_{\max}(0) \ll \text{Ampl.}(0)_{\max} \equiv \frac{8\pi}{\sqrt{3n}} \frac{\Gamma(1/n)}{\Gamma(1/n + 1/2)} \sqrt{\Omega_{OQ}} M_{pl} \sim M_{pl} \end{cases}$$

$$HT \ll 1 \Rightarrow z \ll z_{\max} \equiv F_n^{-1}(1/H_0 T_0) \quad \text{for given } H_0 T_0$$

$z_{\max} \sim$  oscillation starting time

$$F_n(z) \equiv (1+z)^{3(2-n)/(2+n)} \sqrt{\rho/\rho_c}$$

# Oscillating Quintessence

(example)

$$n = 0.1, \quad w_{OQ} = -0.9$$

$$M \gg M_{\min} = 4.3 \times 10^{-4} \text{ eV}, \quad \frac{|\phi|_{\max}(0)}{M_{pl}} \ll \frac{\text{Ampl}(0)_{\max}}{M_{pl}} = 40$$

$M$	$(\rho_c^{1/4})$ $10^{-2.6} \text{ eV}$	eV	GeV	TeV	(GUT) $10^{15} \text{ GeV}$	$(M_{pl})$ $10^{19} \text{ GeV}$
$\text{Log} \left[ \frac{ \phi _{\max}(0)}{M_{pl}} \right]$	-32	-133	-484	-601	-1069	-1225
$\text{Log}[H_0 T_0]$	-30	-131	-482	-599	-1067	-1223
$\text{Log}[z_{\max}]$	6.9	28	103	128	227	260

# Summary

- Feasibility of **OK** oscillation  $\rightarrow$  repulsive gravity  $\rightarrow$  cosmic acceleration  
(dark energy, enough negative pressure)

(opposite to the naive thinking:  
oscillation  $\rightarrow$  excitation particle  $\rightarrow$  attractive gravity  $\xrightarrow{\text{X}}$  cosmic acceleration)

- $w_{\text{OQ}}$  is simply determined by the power  $n$ :

$$w_{\text{OQ}} = \frac{n-2}{n+2}$$

$n$	4	2	1	0.1	$0^+$
$w_{\text{OQ}}$	1/3	0	-1/3	-0.9	$-1^+$

- Requirements from the oscillation condition and the observational constraints are presented.

Extend the scope of DE model construction.

For models involving scalar field(s) such as quintessence (even other fields)